Glueball Searches: Experiment

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Outline:

- Looking for glue: Where and How
- History and puzzles
 - The scalars $f_0(1710)$, $f_0(1370)$, and $f_0(1500)$
 - The pseudoscalars $\eta(1418)$ and $\eta(1475)$
- Resolving the puzzles
 - High statistics $J/\psi \to \gamma X$ using partial waves
 - $J/\psi \to \gamma X$ followed by $X \to \gamma Y$
- $M_X \ge 2 \text{ GeV}/c^2$: Tensor glueballs and J = 4 mesons

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Looking for glue: Where? Morningstar & Peardon arXiv:nucl-th/0309068



Key Point: Masses and quantum numbers look like ordinary $q\bar{q}$ mesons

 $\Rightarrow You need to consider$ the *dynamics* of production and decay.

So, how do we create "glue rich" situations?

Looking for glue: *How?*

p̄p Annihilation
C. Amsler,
Rev.Mod.Phys.
70(1998)1293



pp Central Collisions
F. Close & A. Kirk,
Phys.Lett.B 397(1997)333



J/ψ Radiative Decay
F. Close, G. Farrar & Z. p.Li,
Phys.Rev.D55(1997)5749
Main focus of this talk!



History and Puzzles

More states found than are predicted by the quark model. \Rightarrow Could the extra states be glueballs?

• The scalars $f_0(1710)$, $f_0(1370)$, and $f_0(1500)$

Quark Model predicts only two. (These are the isoscalar $u\bar{u} + d\bar{d} \equiv n\bar{n}$ and $s\bar{s}$.) Prime suspect for the lightest glueball.

• The pseudoscalars $\eta(1418)$ and $\eta(1475)$

Quark Model predicts only one. (This is the $\eta'(958)$ radial excitation). Mass disagrees with lattice QCD.

 $J/\psi
ightarrow \gamma X$ is a key dynamical ingredient!

The Scalars $f_0(1710)$, $f_0(1370)$, and $f_0(1500)$

Prime tool is $J/\psi \to \gamma X$ with $X \to YY$ and spin-zero Y $f_0(1710)$

Early glueball excitement @ SPEAR: The $\theta(1700)$ Xtal Ball $(J/\psi \rightarrow \gamma \eta \eta)$ and Mark-III $(J/\psi \rightarrow \gamma K^+K^-)$ Spin/Parity assignments using Partial Wave Analysis Mark-III (W. Dunwoodie) in SLAC-PUB-7163 (1997) BES: J. Z. Bai *et al.* Phys. Rev. D 68, 052003 (2003)

 $f_0(1370)$

Broad S-wave enhancement known since 1970's. It seems to be clear in $J/\psi \rightarrow \gamma \pi^+ \pi^-$.

 $f_0(1500)$

Excitement in 1990's: Narrow $\pi^0 \pi^0$ state in $\bar{p}p \to 3\pi^0$. Little or no evidence in $J/\psi \to \gamma X$!

Mark-III Spin Parity Assignments



Data Points: A Bin-by-Bin (i.e. Mass Independent) Fit Solid Lines: A Mass Dependent Fit to the Data Points

Recent BES: $J/\psi \rightarrow \gamma K \bar{K}$



• $J/\psi \rightarrow \gamma K \bar{K}$ is dominated by $f_0(1710)$ and $f'_2(1525)$ • $J/\psi \rightarrow \gamma \pi^+ \pi^-$ is dominated by $f_0(1710)$ and $f_2(1270)$, ... and $f_0(1370)$ (See also preliminary BES)

The $f_0(1500)$ is Seen Elsewhere



Important Goal: Identify $J/\psi \rightarrow \gamma f_0(1500)$

Note potential interference with $f'_2(1525)$ \Rightarrow Need careful partial wave analysis

BES Global Fit to $J/\psi \rightarrow \gamma K \bar{K}$ shows slight preference for including $f_0(1500)$ in the fit.

Further evidence from BES in $J/\psi \rightarrow \gamma \sigma \sigma \rightarrow \gamma \pi^+ \pi^- \pi^+ \pi^-$ J. Z. Bai *et al.*, Phys. Lett. B 472, 207 (2000)



We should be able to observe this in CLEO-c and BESIII and to independently measure branching ratios for various $f_0(1500)$, $f_0(1710)$, and $f_0(1370)$ decays.

The Case for $J/\psi \to \gamma \pi^0 \pi^0 \to 5\gamma$

 $J/\psi \rightarrow \gamma \pi^+ \pi^-$ can be done, but it is hard: X. Y. Shen. Mod. Phys. Lett. A 18, 340 (2003)



Better: Two identical particles (e.g. $\pi^0 \pi^0$ or $K_S K_S$)

- Such backgrounds will not be observed.
- Only 0^{++} , 2^{++} , ... are allowed

 $J/\psi
ightarrow 5\gamma$ also allows $J/\psi
ightarrow \gamma\eta\eta$

The Pseudoscalars $\eta(1418)$ and $\eta(1475)$

Originally suggested as the "first glueball": M. S. Chanowitz, Phys. Rev. Lett. 46, 981 (1981)

Observed by Mark-II @ SPEAR in the decay $J/\psi \rightarrow \gamma K_S K^{\pm} \pi^{\mp}$ as a peak with $M(K_S K^{\pm} \pi^{\mp}) = 1440$ MeV.

This became the " E/ι Puzzle". We now understand this region as populated by three different resonances:

 $egin{aligned} f_1(1420) &
ightarrow ar{K}K^\star + ext{c.c.} & ext{Accomodated by Quark Model} \ \eta(1418) &
ightarrow a_0(980)\pi & ext{Only one of these...} \ \eta(1475) &
ightarrow ar{K}K^\star + ext{c.c.} & ext{can be accomodated} \end{aligned}$

Could the extra η be a glueball? No experimental evidence against this, but it would disagree with lattice gauge calculations.

Data on $J/\psi \rightarrow \gamma K \bar{K} \pi$



itz plot for $1.35 < M_{K_S^0 K^{\pm} \pi^{\mp}} < 1.45 \text{ GeV}/c^2$; (c) Dalitz plot for $1.45 < M_{K_S^0 K^{\pm} \pi^{\mp}} < 1.55 \text{ GeV}/c^2$.

However, see also BES results: J. Z. Bai *et al.*, Phys. Lett. B 440, 217 (1998) J. Z. Bai *et al.*, Phys. Lett. B 476, 25 (2000)

 $M_{K_c^{\circ}K^{\pm}\pi^{\mp}}$ (GeV/c²)

Resolving the Puzzles: I High Statistics $J/\psi \to \gamma X$

Partial Wave Analysis relies on subtle differences in the various angular distributions.



 \Rightarrow Need plenty of statistics not only for quality of fit, but also to test sensitivity of acceptance to various cuts.

Note: Good place to start is with a well understood detector with as flat an acceptance function as possible!

Can one interfere $J/\psi \to \gamma K \bar{K}$ with $J/\psi \to K^* \bar{K}$?



 \Rightarrow We might be able to exploit this for higher mass KK resonances *if* there are high enough statistics.

Resolving the Puzzles: II Radiative Decay as a Flavor Probe

F. E. Close, A. Donnachie and Y. S. Kalashnikova, "Radiative decays: A new flavor filter," Phys. Rev. D 67, 074031 (2003)

Fundamental problem: How do the $n\bar{n}$, $s\bar{s}$, and glueball scalars mix to form the $f_0(1710)$, $f_0(1370)$, and $f_0(1500)$?

Hadronic decays can be used, but glueball is hard to model: F. E. Close and A. Kirk, Eur. Phys. J. C 21, 531 (2001)

Photons only couple to the quarks. $\Rightarrow Radiative \ decay \ width \ should \ be \ very \ sensitive \ to \ the \ glueball/q\bar{q} \ admixture.$

Results of Model Calculations

Three Scenarios, based on values of bare masses:

 $\begin{array}{ll} M(glueball) < M(n\bar{n}) & \equiv ``L" \\ M(n\bar{n}) < M(glueball) < M(s\bar{s}) & \equiv ``M" \\ M(glueball) > M(s\bar{s}) & \equiv ``H" \end{array}$

Radiative Decay Widths in keV							Γ_{Tot}
	f_0 -	$ ightarrow \gamma ho('$	770)	f_0 -	$ ightarrow \gamma \phi$	(1020)	MeV
State	L	\mathbf{M}	Η	L	\mathbf{M}	Η	
$f_0(1370)$	443	1121	1540	8	9	32	~ 300
$f_0(1500)$	2519	1458	476	9	60	454	109
$f_0(1710)$	42	94	705	800	718	78	125

Excellent discrimination!

 $\Rightarrow \text{Expect } \mathcal{B}(f_0 \to \gamma V) \approx 10^{-2} \text{ to } 10^{-4}.$ For $\mathcal{B}(J/\psi \to \gamma f_0) \approx 10^{-3}$ we should acquire 10,000 to 100 events for $10^9 J/\psi$.

$M_X \ge 2 \text{ GeV}/c^2$: Tensor glueballs and J = 4 mesons

The $J/\psi \to \gamma X$ mass spectrum shows a lot of structure for $M_X \ge 2 \ {
m GeV}/c^2.$





Quark model predicts many states with $J \leq 4$. Lattice QCD predicts glueballs with $J^{PC} = 2^{++}, 0^{++}, 0^{-+}, 2^{-+}, 1^{+-}.$

Many interferences will require high statistics and excellent control of systematics to disentangle the states.

Conclusions

Radiative transitions between vector and scalar mesons will provide strong constraints on the presence and structure of glueballs with $1.3 \leq M \leq 2.0 \text{ GeV}/c^2$.

High statistics $J/\psi \to \gamma X$, including $X \to \gamma \{\rho, \phi\}$, may be within reach of CLEO-c and BESIII.

A complete analysis will be limited by data volume. Will likely need $\geq 10^9 J/\psi$'s.

The region with $M_X \ge 2 \text{ GeV}/c^2$ needs to be carefully explored for narrow and/or broad resonances. Partial Wave Analysis will be a necessary tool.

High statistics are good, but excellent knowledge of the detector acceptance is also crucial.