Precision Charm Experiment and Precision LQCD

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Precision Lattice QCD ?

BESIII/CLEO-c promise high precision measurements of D meson decay constants, form factors, *etc.* — typically 2% accuracy.

The impact on/from lattice QCD is

- provide calibration for lattice calculation of those quantities.
- give precise determination of some CKM matrix elements, provided that lattice calculation is available with comparable accuracy.

High precision lattice QCD means a few % accuracy. The CLEO-c report assumed 1%.

An example

Leptonic decay:

$$\Gamma(D^+ \to \mu^+ \nu_\mu) \;\propto\; |V_{cd}|^2 f_D^2$$

- assuming the CKM unitarity
  $\Rightarrow$ determination of f_D: 2.3%.
  $\Rightarrow$ calibration of the lattice calculation
- relying on the lattice calculation ($\sim$ 1%)
  $\Rightarrow$ determination of $|V_{cd}|$: 2.3%.

Model independent lattice calculation is as important as the precise experiments.

Other examples

exp + lattice $\Rightarrow$ CKM elements

- leptonic decay constants: $D_{(s)} \to l\nu$
- semileptonic form factors: $D \to Kl\nu$, $\pi l\nu$

exp + lattice $\Rightarrow$ deeper understanting of QCD

- quarkonium spectrum
- glueballs, hybrids, other exotics ($\to$ Morningstar)

I will mainly consider the first class.

Specific questions (to myself)

- Is the 1% (or a few %) accuracy really achievable in the next several years?
- It must include the effect of dynamical quarks (up, down, and strange). Is it feasible?
- What is needed to achieve this goal?

To consider these questions, let us look back what happened in the past 10 years.

Plan of the talk

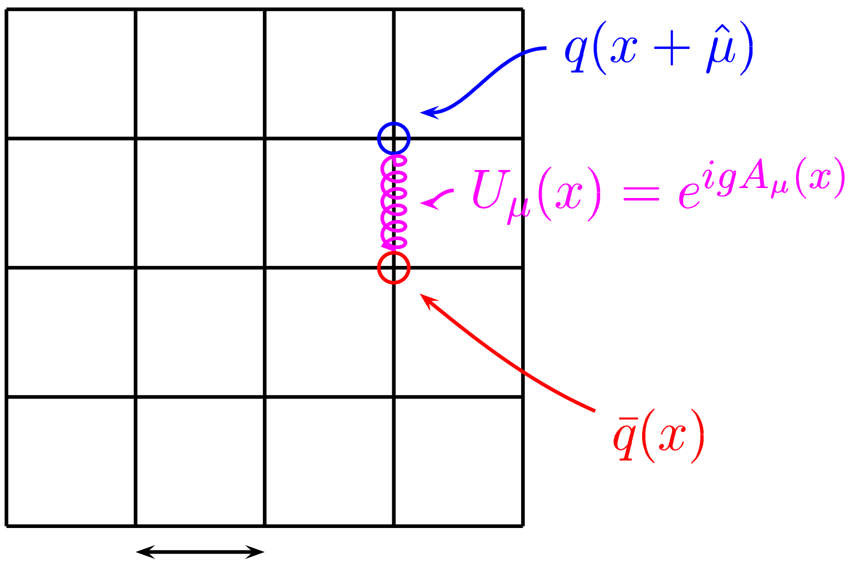
1. Improvements in lattice QCD
   - Symanzik's improvement
   - HQET/NRQCD
   - renormalized perturbation theory
2. Unquenching
   - why so hard?
   - chiral extrapolation
   - fermion actions
3. Future — Is the 1% feasible?
   - machines

Improvements in lattice QCD

- Introduction
- Symanzik's improvement
- HQET/NRQCD or conventional
- renormalized perturbation theory
  (or non-perturbation matching)

Lattice QCD = first principles calculation

A regularization of QCD:



lattice spacing a

- Numerical simulation is possible.
  path integral $\Rightarrow$ Monte Carlo
- It gives a nonperturbative formulation of QCD.
  $\Leftrightarrow$ Dimensional regularization is defined through perturbation theory.

prediction of LQCD = prediction of QCD

Ideally ...

To reproduce the real world, one needs

- unquenched, N_f = 2+1.
- L = 5 fm.
- a = 0.02 fm;
  or a^{-1} = 10 GeV.
- m_{ud} = several MeV,
  m_s = 100 MeV.
- statistics $\sim$ 10K.

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Empirial law : the computational demand scales as

$$\left[\frac{m_\pi/m_\rho}{0.6}\right]^{-6}\left[\frac{L}{3\ \mathrm{fm}}\right]^{5}\left[\frac{a^{-1}}{2\ \mathrm{GeV}}\right]^{7}$$

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For this example, we need

10^{10} TFlops · year

Theoretical/algorithmic improvements are crucial.

Role of effective theories

It is hard to describe the physics at different energy scales on a single lattice.

$$m_W \gg m_b > m_c \gg \Lambda_{QCD} \gg m_q$$

Lattice QCD deals with the physics at the $O(\Lambda_{QCD})$, leaving the others for effective theories.

- m_W : Weak effective Hamiltonian (4-fermion interactions)
- m_b, m_c : Heavy Quark Effective Theory (HQET)
- m_q : Chiral Perturbation Theory (ChPT)

- $1/a$: Symanzik's effective theory (discretization error)

Symanzik's effective theory

How the discretization error looks like:

$$\mathcal{L}_{lat} \doteq \mathcal{L}_{QCD} + \mathcal{L}_I$$

$\doteq$ means "give the same on-shell amplitude."
$\mathcal{L}_{QCD}$ is the continuum QCD lagrangian.

- Discretization error is described by local operators $\mathcal{L}_I$:

$$\mathcal{L}_I = a\mathcal{L}^{(5)} + a^2\mathcal{L}^{(6)} + \cdots$$

- Theoretical basis to construct the *improved* actions.

Improved actions

Order counting assuming $\Lambda_{\mathrm{QCD}} = 400$ MeV:

a (fm)	0.2	0.1	0.05
$1/a$ (GeV)	1	2	4
$O(a\Lambda_{\mathrm{QCD}})$	40%	20%	10%
$O((a\Lambda_{\mathrm{QCD}})^2)$	16%	4%	1%
$O((a\Lambda_{\mathrm{QCD}})^3)$	6%	1%	$<$ 1%
$O((a\Lambda_{\mathrm{QCD}})^4)$	3%	$<$ 1%	$<$ 1%

To achieve the 1% accuracy,

- $O(a)$-improved action + extrapolation in a^2
- $O(a^2)$-improved action at a = 0.1 fm.

HQET

Heavy quark is static (non-relativistic) in the heavy-light (heavy-heavy) meson. Dynamical degrees of freedom are $\sim O(\Lambda_{\mathrm{QCD}})$, which we treat non-perturbatively on the lattice.

$$\mathcal{L}_Q = Q^\dagger \left[iD_0 + \frac{\boldsymbol{D}^2}{2m_Q} + \cdots \right] Q$$

$\cdots$ an expansion in $\Lambda_{\mathrm{QCD}}/m_Q$.

- HQET: Eichten *et al.* (1990)
- NRQCD: Lepage *et al.* (1992)
- Fermilab action: El-Khadra, Kronfeld, Mackenzie (1997)

HQET order counting

Assuming $\Lambda_{\rm QCD} = 400$ MeV:

m_Q (GeV)	m_b 4.5	m_c 1.5
$O(\Lambda_{\rm QCD}/m_Q)$	9%	27%
$O((\Lambda_{\rm QCD}/m_Q)^2)$	1%	7%
$O((\Lambda_{\rm QCD}/m_Q)^3)$	$<$ 1%	2%
$O((\Lambda_{\rm QCD}/m_Q)^4)$	$<$ 1%	$<$ 1%

To achieve the 1% accuracy,

- $O(\Lambda_{\rm QCD}/m_Q)$ or $O((\Lambda_{\rm QCD}/m_Q)^2)$ action for b quark
- $O((\Lambda_{\rm QCD}/m_Q)^3)$ action for c quark

Without HQET

It is also possible to use the conventional fermion action as far as $am_Q \ll 1$. For m_c = 1.5 GeV,

a (fm)	0.1	0.05	0.03	0.02
$1/a$ (GeV)	2	4	6.7	10
$O(am_c)$	75%	38%	22%	15%
$O((am_c)^2)$	56%	14%	5%	2%
$O((am_c)^3)$	42%	5%	1%	$<$ 1%

To achieve the 1% accuracy,

- $O(a)$-improved action at $a \lesssim$ 0.03 fm + extrapolation in a^2.
- $O(a^2)$-improved action at $a =$ 0.03 fm.

Example: Continuum extrapolation of f_{D_s}

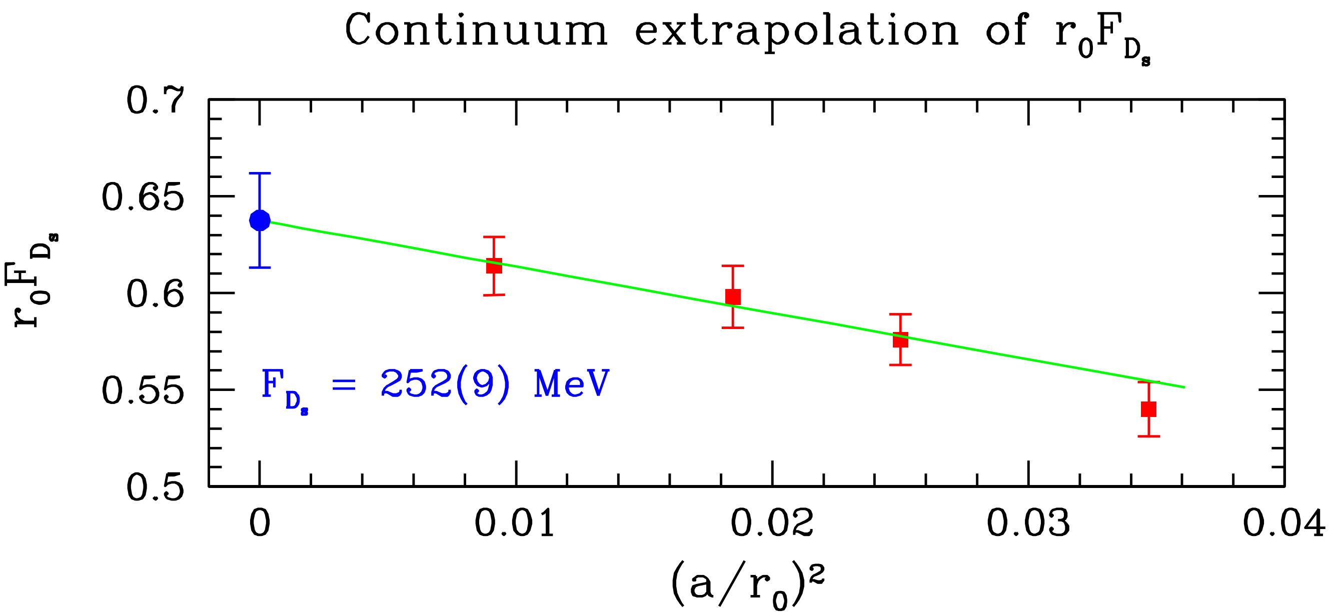
Juttner, Rolf, PLB560 (2003) 59.

- quenched approximation
- $O(a)$-improved action
- $a =$ 0.09–0.05 fm

Extrapolation in a^2:

$$f_{D_s} = 252(9)\ \text{MeV}.$$

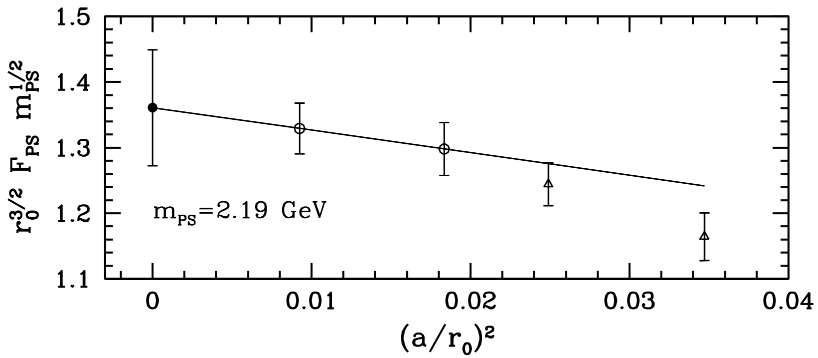
(4% error) using f_K as an input for the lattice scale.

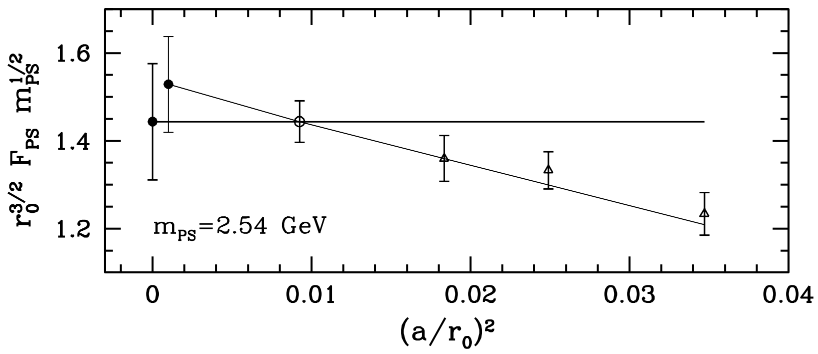


Extrapolation to the b quark

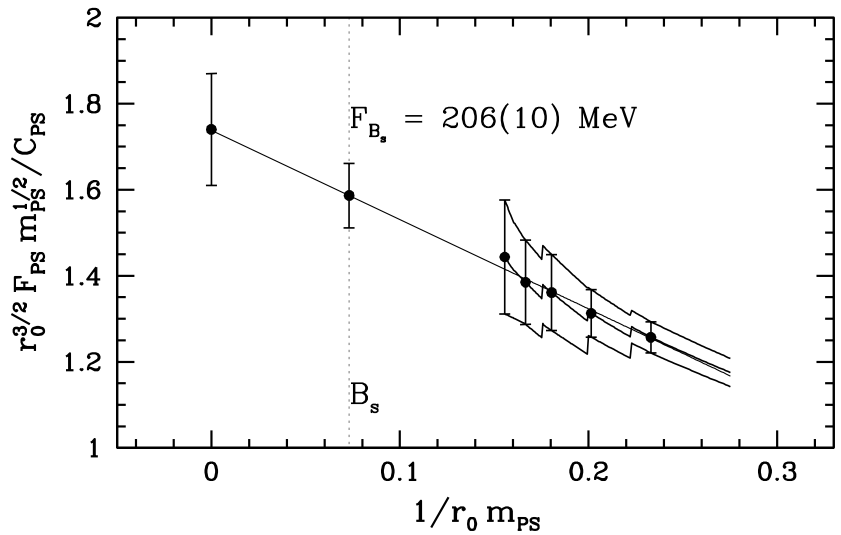
Rolf *et al.*, hep-lat/0309072

Continuum extrapolation at several m_Q,





Then, another extrapolation (or interpolation) in $1/m_Q$.



Better controled with a combination with HQET.

Perturbative matching

Matching of continuum and lattice operators

$$\mathcal{O}^{\overline{\rm MS}}(\mu) = Z(\mu a)\mathcal{O}^{lat}(a)$$

In most cases, $Z(\mu a)$ is known only at one-loop.

Renormalized lattice perturbation theory

Lepage, Mackenzie (1993),

$$1 + c_1\alpha_0 + c_2\alpha_0^2 + \cdots$$ poor convergence

$$= \ 1 + c_1 \boxed{\alpha_V(q^*)} + c_2'\alpha_V^2(q^*) + \cdots$$ much better

renormalized coupling

Perturbative error

Coupling constant is evaluated at a typical scale $q^\star \sim 2/a$.

a (fm)	0.2	0.1	0.05
$1/a$ (GeV)	1	2	4
$\alpha_s(2/a)$	0.26	0.22	0.18
$O(\alpha_s)$	26%	22%	18%
$O(\alpha_s^2)$	7%	5%	3%
$O(\alpha_s^3)$	2%	1%	$<$ 1%

To achieve the 1% accuracy,

- two-loop calculation at $a \lesssim$ 0.1 fm; need automated perturbative calculation.

Non-perturbative matching

Or, one may prefer some non-perturbative methods to eliminate the perturbative error.

Heitger, Sommer, hep-lat/0310035.

- Matching the relativistic lattice action and HQET for $am_Q \ll 1$. It is possible if the entire lattice volume is small $L_0 \simeq$ 0.2 fm.
- Recursively match the HQET in larger volumes $L_i = 2^i L_0$ until L_i becomes physical volume 2 fm.

Both the perturbative and non-perturbative avenues are to be pursued.

Unquenching

- why so hard?
- chiral extrapolation

Quenched approximation

An “approximation” to neglect the fermion determinant in the Feynman path integral,

$$Z = \int [dU_\mu] \, \det M^2 \, e^{-S_G}$$

due to its huge computational demand.

- Most lattice calculations ($\lesssim$ 2000) were within the quenched approximation.
- Its uncertainty is out of control. The only possoble solution is to put $\det$ back.

Why so hard?

Monte Carlo simulation has to deal with

$$Z = \int [dU_\mu] \, \det M^2 \, e^{-S_G} = \int [dU_\mu][d\phi] \, e^{-S_G - \phi^\dagger (M^\dagger M)^{-1} \phi}$$

$M \equiv \det(\not{D}[U_\mu] + m)$ is the fermion matrix; ϕ is a (fictitious) pseudo-fermion field.

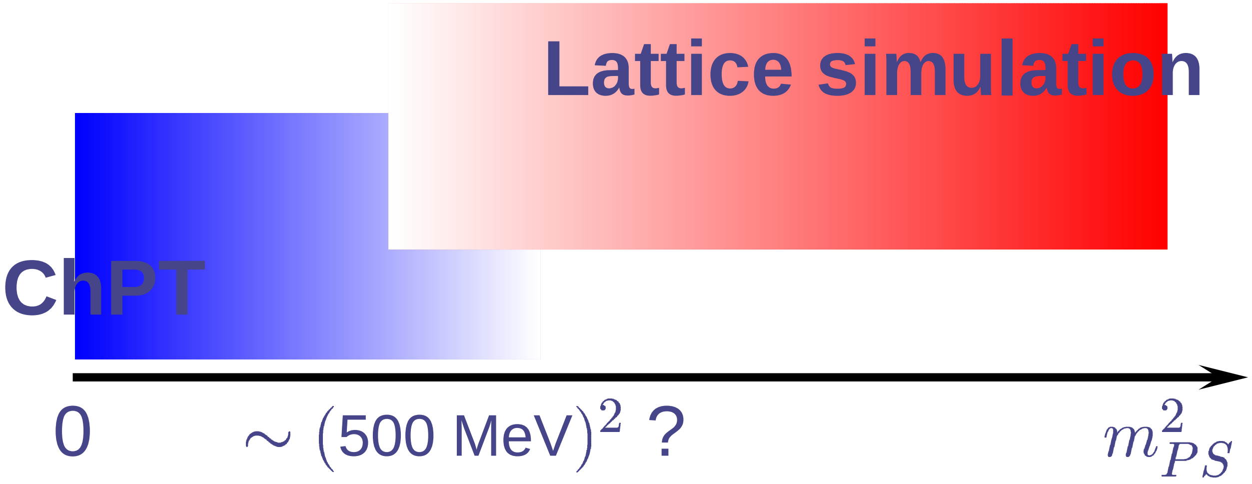
- The effective action becomes non-local $\phi^\dagger (M^\dagger M)^{-1} \phi$; local updation is difficult.
- Matrix inversion $(M^\dagger M)^{-1}$ is time-consuming especially for light quarks.

Simulation of light dynamical quarks is very hard: $\sim 1/m_q^3$.

How small sea quark masses needed?

QCD with very small quark masses is described by Chiral Perturbation Theory (ChPT)

- The chiral extrapolation of lattice data must be consistent with ChPT near $m_q = 0$.
- Test whether the observed quark mass dependence is consistent with the ChPT formula (especially, the chiral log).



Test for the pion decay constant

Full ChPT

$$\frac{f_{SS}}{f} = 1 - \frac{N_f}{2} y_{SS} \ln y_{SS} + \frac{1}{2}\alpha_5 y_{SS} + \frac{1}{2}\alpha_4 N y_{SS}$$

$$y_{SS} \equiv \frac{m_{SS}^2}{(4\pi f)^2}$$

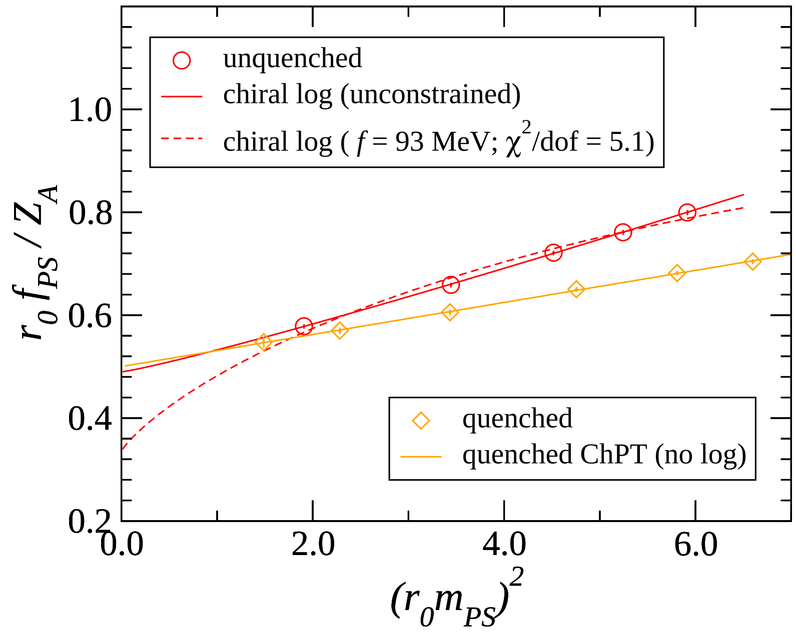
—chiral log with a known coeffi cient

Quenched ChPT

$$\frac{f_{11}}{f} = 1 + y_{11}\frac{1}{2}\alpha_5.$$

—no quenched chiral log

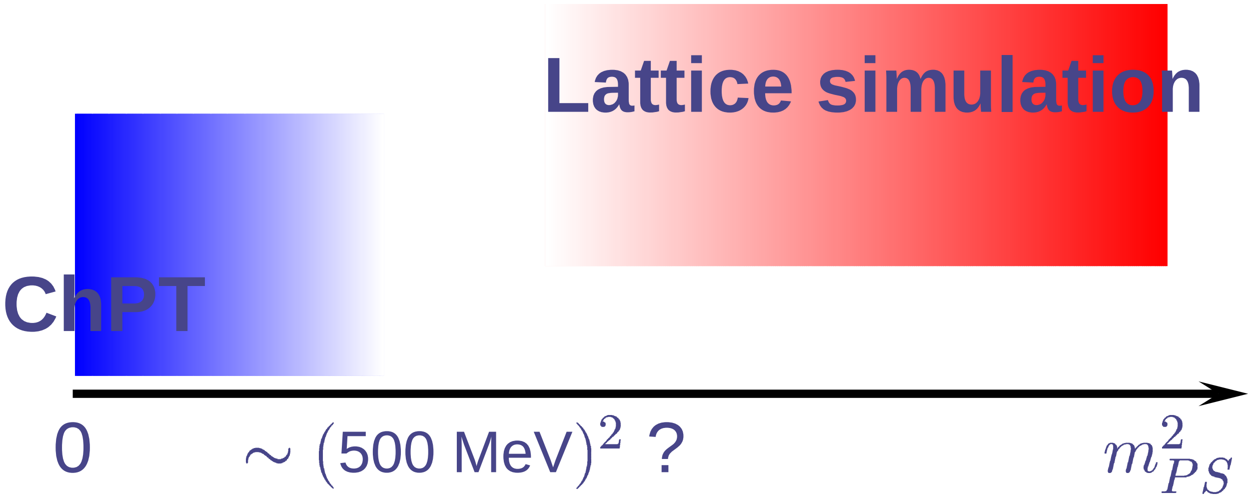
JLQCD (2002): a high statistics test with the $O(a)$-improved Wilson fermion.



Data do not show the curvature characteristic of the chiral log.

Possible interpretation

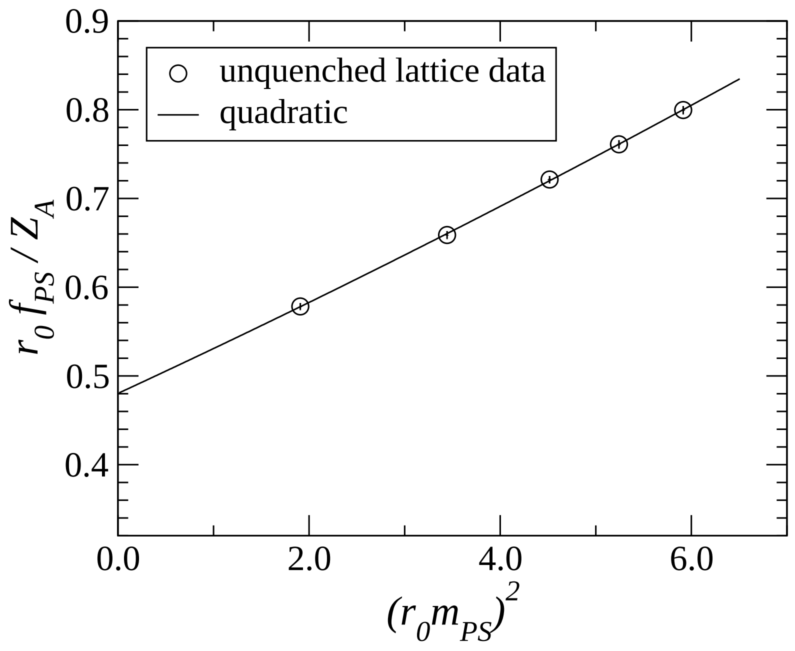
Lattice data lie beyond the reach of ChPT.



- Need much smaller sea quark masses, probably corresponding to $m_{PS} \simeq$ 300 MeV.
- Otherwise, the chiral extrapolation introduces significant uncertainty.

Impact on physical quantities

decay constant f_π



Possible fit forms:

- quadratic fit (no chiral log)

Impact on physical quantities

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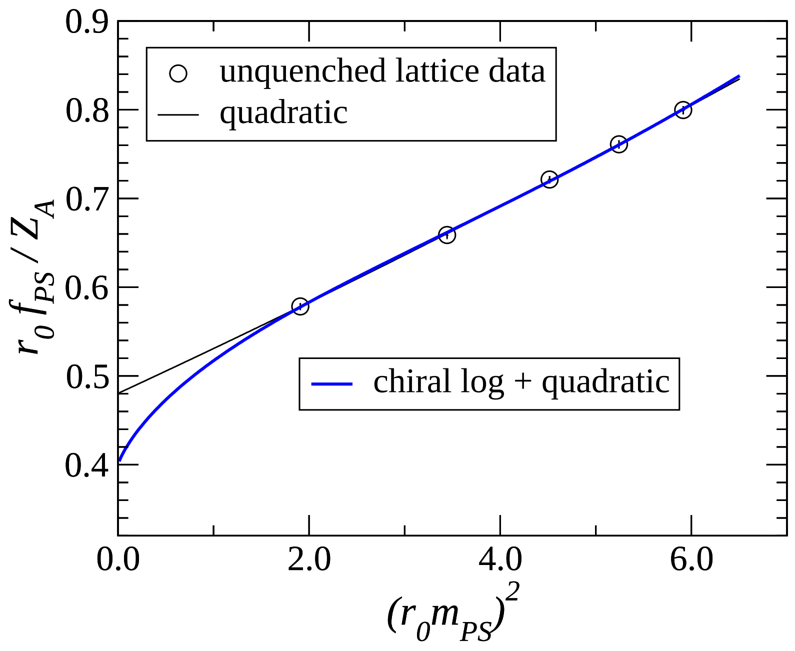

Possible fit forms:

- quadratic fit (no chiral log)
- chiral log (with the known coefficient) plus quadratic

  curvature cancels in the data region.

Impact on physical quantities

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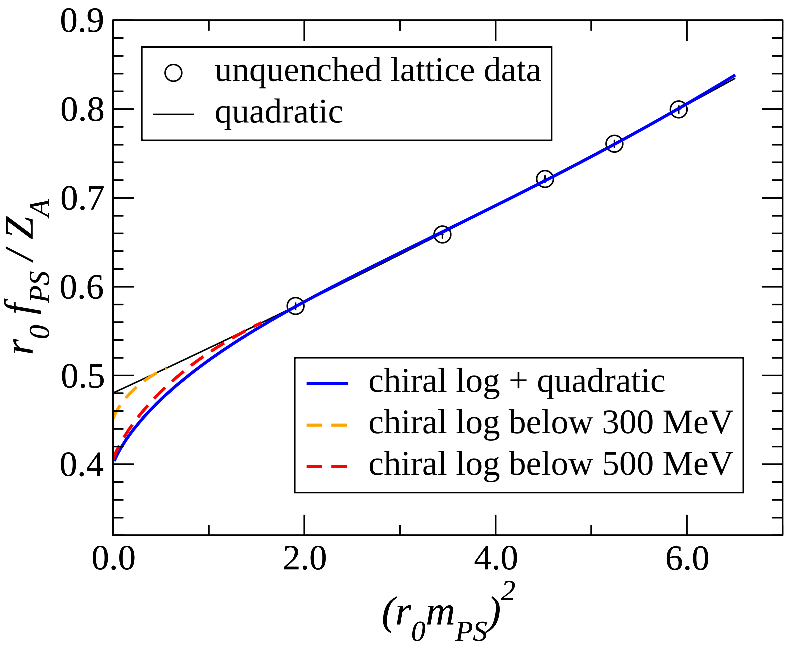


Possible fit forms:

- quadratic fit (no chiral log)
- chiral log (with the known coefficient) plus quadratic

  curvature cancels in the data region.
- One-loop ChPT formula below μ (μ = 300 MeV and 500 MeV are shown.)

Uncertainty of order $\pm$10% in the chiral limit.

Heavy-light meson decay constant

ChPT for the heavy-light decay constant (N_f = 2)

Grinstein *et al.* (1992)

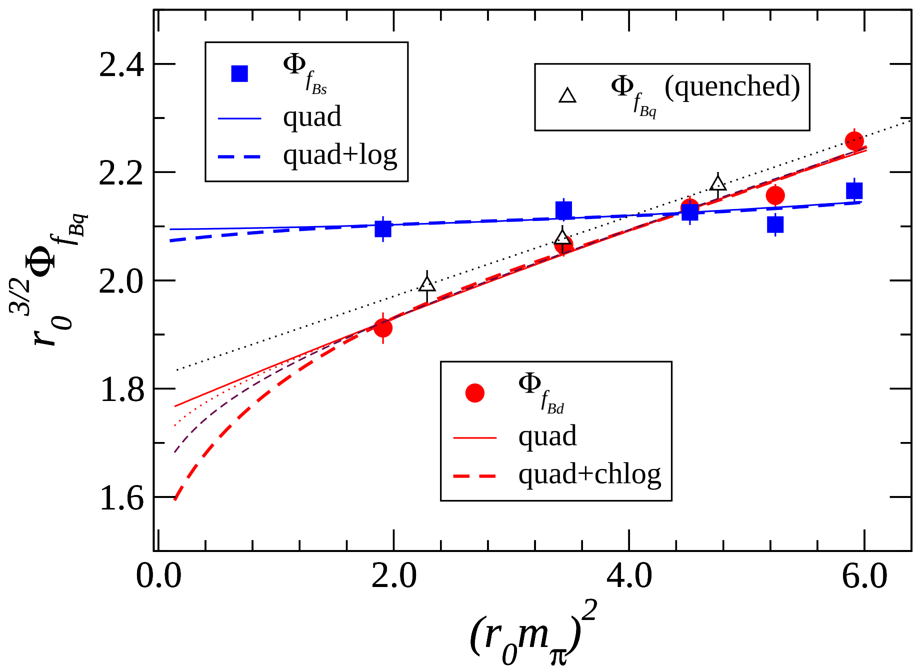
$$\frac{f_B}{f_B^{(0)}} = 1 - \frac{3}{8}(1+3g^2)\, y_{SS} \ln y_{SS}$$
$$+ \text{analytic terms}$$
$$y_{SS} \equiv \frac{m_{SS}^2}{(4\pi f)^2}$$

in the heavy quark mass limit.

g : $B^* B\pi$ coupling

$$g = 0.59(7) \quad (D^*\text{width}, \text{CLEO})$$

JLQCD (2003)



Significant uncertainty depending on the form of chiral extrapolation.

Chiral symmetry on the lattice

The problem of chiral extrapolation may be related to the problem of chiral symmetry on the lattice.

- Wilson-type fermions:
  - Add a Wilson term $\frac{1}{2}\bar{\psi}D^2\psi$ to the action.
  - A conventional choice in the quenched calculations.
  - Chiral symmetry is explicitly broken; massless limit is not determined by any symmetry.
    $\Rightarrow$ The computational time $1/(m_q + \Delta m_q)^n$ fluctuates, or even diverges, configuration by configuration.
  - Lightest available sea quark mass is $\sim m_s/2$.

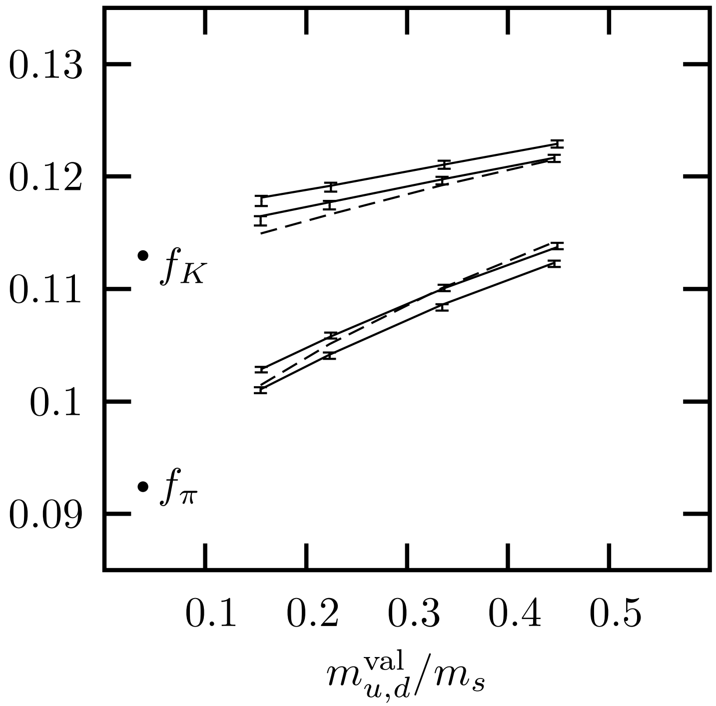
- Staggered fermions:
  - Contains 4 flavors (4 *tastes* in modern terminology) of quarks; Chiral U(1) remains out of SU(4).
  - Complicated flavor (taste) structure: 15 pions (1 is Nambu-Goldstone of U(1); others are not), 64 protons *etc.*
  - Take $(\det M)^{1/4}$ per flavor.
    $\Rightarrow$ Effective action is non-local; inconsistent as a lattice field theory.
  - Effective lattice spacing is $\times 2$ larger.
  - Numerically so cheap. Lightest available sea quark mass is $\sim m_s/6$.

Realistic staggered simulations

HPQCD-UKQCD-MILC-Fermilab collaboration (2003)

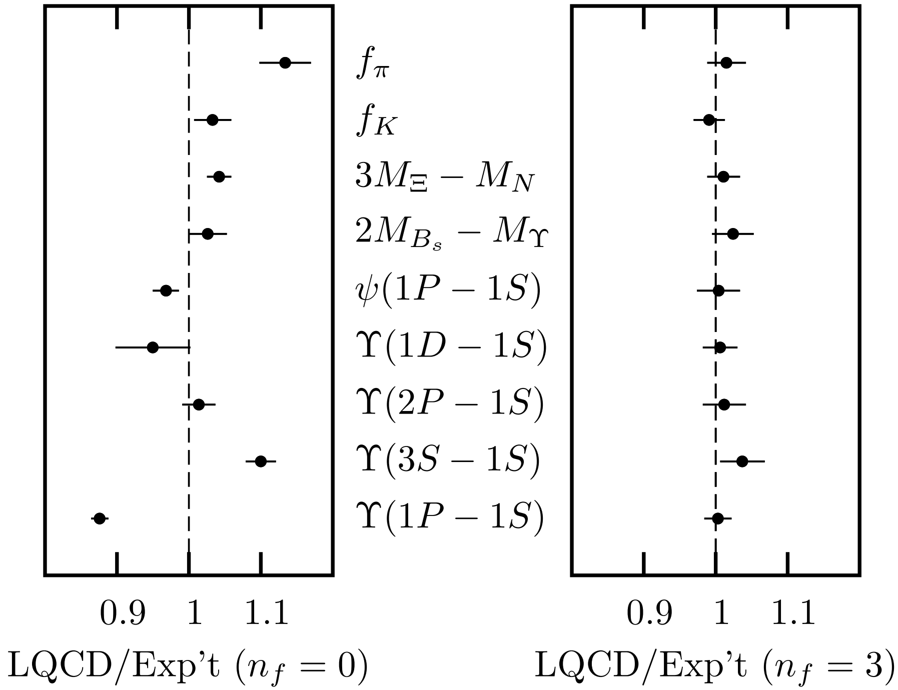
- 2+1 flavor (u, d and s)
- $O(a^2)$-improved staggered fermion
- $m_{u,d} = m_s/2 - m_s/6$
  chiral extrapolation is done with the data below $m_s/2$.
- lattice spacing 1/8 fm and 1/11 fm.

decay constant



First results

quenched versus unquenched (hep-lat/0304004).



Very impressive agreement with experiments.
Promising also for B and D physics.

Other choices

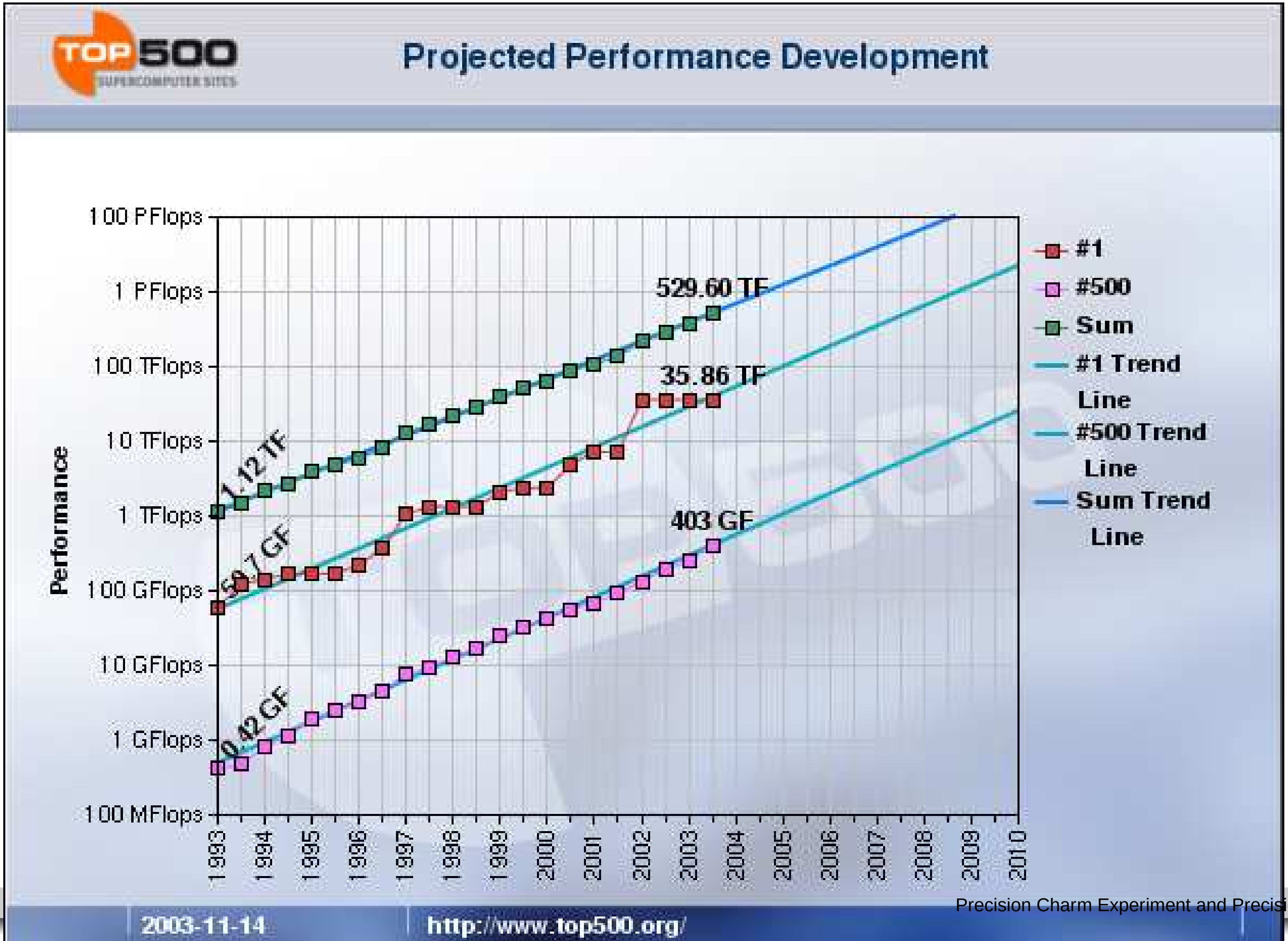
- The best available fermion formulation is the Ginsparg-Wilson fermions (domain-wall or overlap).
  - An exact chiral symmetry on the lattice without introducing fictitious tastes.
  - Tested on the quenched lattices. Simulations with very light quark masses are possible.
  - The unquenched simulation is extremely demanding (a factor $\times$ 10–100 over the Wilson-type).

Future

– machines
– perspectives

Machine

Moore's law: The computer speed becomes $\times 10$ in 5 years.



Machines available for lattice QCD

Not a complete list:

- SR8000 (KEK): 1.2 TFlops since 2000.
- QCDSP (Columbia + RIKEN-BNL): 1.1 TFlops (total) since 1998.
- QCDOC: 5–10 TFlops in 2004 @ Columbia, RIKEN-BNL, Edinburgh.
- apeNEXT: more than 5 TFlops in 2004 (?) @ INFN, DESY, ...
- many PC clusters @ many places

In 2005–2010, several tens of TFlops will be available for lattice calculations.

Perspectives (1)

I don't see any fundamental problems to achieve the goal, *i.e.* the 1% accuracy for charm physics.

- $O(a^2)$-improved action at $a = 0.1$ fm.
- $O((\Lambda_{\mathrm{QCD}}/m_Q)^3)$ action for c quark. (Without HQET, we need the $O(a^2)$-improved action at $a = 0.03$ fm.)
- two-loop matching at $a \lesssim 0.1$ fm.

All these iterms are within reach. Actually, they are on the program of the HPQCD-UKQCD-MILC-Fermilab group.

This argument is based on an order counting. Scaling test will be needed to convince ourselves.

Perspectives (2)

Computationally most demanding item is the

- dynamical fermion at $m_{u,d} \lesssim m_s/3$.

At present (with $O(1)$ TFlops machines), this is only feasible with the (improved) staggered fermion. Other fermion formulations will need at least $\times$ 10 more computer time = five more years to follow.

Therefore,

- *short–mid term:* More test of the improved staggered fermion (scaling, taste breaking, *etc.*)
- *mid–logn term:* Need much faster algorithms for other fermion formulations (especially the Gisparg-Wilson fermions)

Perspectives (3)

Then, as stated in the CLEO-c report, the impact of the charm factories is

- To determine the CKM elements $|V_{cs}|$, $|V_{cd}|$ with a few % accuracy.
- To give calibration data for lattice QCD calculation, and thus to help the B physics program, *i.e.* determination of $|V_{ub}|$, $|V_{td}|$, $|V_{ts}|$, and other form factors.

An integral part of the flavor physics.